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### Short communication

# Ionic conduction in $Sn_{1-x}Sc_xP_2O_7$ for intermediate temperature fuel cells

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#### ABSTRACT

A novel series of mixed ion conductors,  $Sn_{1-x}Sc_xP_2O_7$  (x = 0.03, 0.06, 0.09, 0.12), were synthesized by a solid-state reaction method. The conduction behaviors of the ion conductors in wet hydrogen atmosphere were investigated by some electrochemical methods including AC impedance spectroscopy, gas concentration cells in the temperature range of 323–523 K. It was found that the doping limit of  $Sc^{3+}$  in  $SnP_2O_7$  was between 9 mol% and 12 mol%. The highest conductivity was observed to be  $2.76 \times 10^{-2}$  S cm<sup>-1</sup> for the sample of x = 0.06 under wet H<sub>2</sub> atmosphere at 473 K. The ionic conduction was contributed mainly to proton and partially to oxide ion in wet hydrogen atmosphere from 373 K to 523 K. The H<sub>2</sub>/air fuel cells using  $Sn_{1-x}Sc_xP_2O_7$  (x = 0.03, 0.06, 0.09) as electrolytes (1.7 mm in thickness) generated the maximum power densities of 11.16 mW cm<sup>-2</sup> for x = 0.03, 25.02 mW cm<sup>-2</sup> for x = 0.06 and 14.34 mW cm<sup>-2</sup> for x = 0.09 at 423 K, respectively. The results indicated that  $Sn_{1-x}Sc_xP_2O_7$  is a promising solid electrolyte system for intermediate temperature fuel cells.

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#### 1. Introduction

High temperature protonic conduction in perovskite type oxides such as doped ACeO<sub>3</sub> (A=Sr, Ba), AZrO<sub>3</sub> (A=Ca, Sr, Ba), etc. has received much attention for their potential applications in some electrochemical devices and membrane reactors [1–5]. However, high sintering and operating temperatures for these materials not only lead to metal evaporation but also bring some difficulties in choosing sealing and bonding materials. Therefore, there is a great interest in exploring new electrolyte materials with desired performances at lower sintering and operating temperatures. Considerable effort has been devoted toward the intermediate temperature fuel cells (ITFCs) based on CsHSO<sub>4</sub>, CsH<sub>2</sub>PO<sub>4</sub>, etc. in the temperature range of 423–673 K [6,7]. Recently, numbers of tetravalent metal pyrophosphates AP<sub>2</sub>O<sub>7</sub> (A=Zr, Sn, Ti, Ce, Si, etc.) were reported to have high proton conductivities [8-12]. Nagao et al. reported that the conductivities of  $AP_2O_7$  (A=Sn, Ti, Ge, Si) increased in the order of  $Ge^{4+} < Si^{4+} < Ti^{4+} < Sn^{4+}$  in unhumidified air [9], and that proton-conducting  $Sn_{1-x}R_xP_2O_7$  (R=Al<sup>3+</sup>,  $In^{3+}$ ) were also applied to some electrochemical reactors, such as  $H_2/air$  fuel cell, selective catalytic reduction of NO<sub>x</sub>, direct oxidation of methane to methanol, and gas sensors, etc. [13-18]. Chen et al. reported that Sn<sub>0.9</sub>In<sub>0.1</sub>P<sub>2</sub>O<sub>7</sub> can conduct both protons and

oxide-ions in the temperature range from 403 to 503 K [19]. Protonic conduction in  $Sn_{1-x}R_xP_2O_7$  (R = Sb<sup>3+</sup>, Mg<sup>2+</sup>) was also reported [20,21]. Recently, we reported the protonic and oxide ionic conduction in  $Sn_{1-x}Ga_xP_2O_7$  [22]. However, there is little information about ionic conduction in Sc<sup>3+</sup> doped SnP<sub>2</sub>O<sub>7</sub>. The effective ionic radius (0.0745 nm [23]) of Sc<sup>3+</sup> is close to that (0.069 nm [23]) of Sn<sup>4+</sup> in 6-fold coordination. Therefore, it is interesting to investigate Sc<sup>3+</sup> doped SnP<sub>2</sub>O<sub>7</sub>.

In this study, a novel series of intermediate temperature solid ion conductors,  $Sn_{1-x}Sc_xP_2O_7$  (x = 0.03, 0.06, 0.09, 0.12), were prepared and the conduction behaviors were investigated using some electrochemical methods in the temperature range of 323–523 K. We also made estimation of the protonic and oxide ionic transport properties from the concentration cells of water vapor and hydrogen [19,24,25]. The performances of H<sub>2</sub>/air fuel cell were tested using  $Sn_{1-x}Sc_xP_2O_7$  (x = 0.03, 0.06, 0.09) as electrolytes.

#### 2. Experimental

The initial molar ratios of phosphorus vs. metal ions,  $P_{\rm ini}/({\rm Sn+Sc})$ , should be controlled to be 2.8 due to a fraction of phosphorus loss in the process by vaporization [13,20,22].  ${\rm Sn}_{1-x}{\rm Sc}_x{\rm P}_2{\rm O}_7$  (x = 0.03, 0.06, 0.09, 0.12) was prepared as follows. The required amounts of  ${\rm SnO}_2$  (Luoyang Ship Material Research Institute, 50–70 nm),  ${\rm Sc}_2{\rm O}_3$  (99.9%) and 85% H<sub>3</sub>PO<sub>4</sub> were fully mixed and held with stirring around 623 K until they became solid mixtures. The solid mixtures were ground into powder and

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**Fig. 1.** XRD patterns of the  $Sn_{1-x}Sc_xP_2O_7$  (*x* = 0.03 – 0.12) samples.

heat-treated at 623 K for 30 min until they became white powder, followed by sieving (100 mesh). The resulting powder was pressed into pellets at hydrostatic pressure of 300 MPa, and heattreated at 873 K for 2 h to obtain  $Sn_{1-x}Sc_xP_2O_7$  samples. The crystalline structure of the  $Sn_{1-x}Sc_xP_2O_7$  samples was determined at room temperature by X-ray diffraction (XRD) analysis. The microstructures of the samples were observed by a scanning electron microscope (SEM). Platinum paste was coated on both sides (area: 0.5 cm<sup>2</sup>) of the heat-treated pellets (diameter 20 mm, thickness 1.4-1.8 mm) and heated at 523 K for 30 min. The pellets were connected to Pt wires and sealed between two alumina tubes, and then placed into an electric furnace for electrochemical measurements. The conductivities were measured in wet H<sub>2</sub> atmosphere at 323-523 K by an AC impedance method using electrochemical workstations (Zahner IM6ex). The frequency range for the measurements was 1 Hz to 3 MHz and the ac amplitude was 10 mV. Wet gases were obtained by saturating water vapor at 298 K.

In order to study the ionic conduction in the samples, hydrogen and water vapor concentration cells were constructed and the electromotive forces (EMFs) were measured at 373–523 K. The H<sub>2</sub>/air fuel cell using  $Sn_{1-x}Sc_xP_2O_7$  (x=0.03, 0.06, 0.09) as electrolytes and porous platinum as cathode and anode was constructed. The current–voltage curves were measured to evaluate the performance of the fuel cell.

#### 3. Results and discussion

#### 3.1. X-ray diffraction

The structure of SnP<sub>2</sub>O<sub>7</sub>-based compounds is very complicated. A pseudo-cubic  $3 \times 3 \times 3$  superlattice was observed at room temperature in SnP<sub>2</sub>O<sub>7</sub> by electron diffraction [26]. After heat-treated at 873 K for 2 h, some weak peaks, which cannot be indexed by the cubic structure appeared, but may be indexed by the orthorhombic superlattice, were observed in Sn<sub>0.9</sub>In<sub>0.1</sub>(P<sub>2</sub>O<sub>7</sub>)<sub>1- $\delta$ </sub> prepared by aqueous solution method [10]. Fig. 1 shows the powder X-ray diffraction (XRD) patterns of Sn<sub>1-x</sub>Sc<sub>x</sub>P<sub>2</sub>O<sub>7</sub> samples with various Sc<sup>3+</sup> doping levels. It can be seen from Fig. 1 that the sample of x = 0.03 shows a single cubic phase structure in agreement with that of SnP<sub>2</sub>O<sub>7</sub> in JCPDS (00-029-1352), while the samples of  $x \ge 0.06$  show a weak peak at  $2\theta \approx 23^{\circ}$ , which cannot be indexed by the



**Fig. 2.** SEM images of external surface of (a) x = 0.06 and (b) x = 0.12 pellet samples heat-treated at 873 K for 2 h.

cubic structure, probably due to the formation of the pseudo-cubic  $3 \times 3 \times 3$  superlattice [10]. The lattice parameters for the samples of x = 0.03, 0.06, 0.09, 0.12 were calculated to be a = 0.7941(5) nm. 0.7939(5)nm, 0.7942(4)nm and 0.7935(5)nm, correspondingly. The effective ionic radius (0.0745 nm [23]) of Sc<sup>3+</sup> is larger than that (0.069 nm [23]) of Sn<sup>4+</sup> in 6-fold coordination. Therefore, the lattice parameter is expected to increase with increasing doping level. The sample of x = 0.09 shows the maximum lattice parameter, whereas the sample of x = 0.12 shows the smallest lattice parameter among the determined samples. The significant decrease in lattice parameter of x = 0.12 sample is likely due to the loss of larger scandium in the pyrophosphate phase with the formation of Sc(PO<sub>3</sub>)<sub>3</sub> [10]. This may be also confirmed by comparing the intensities of weak peaks at  $2\theta \approx 23^{\circ}$  of the samples. A weak peak of Sc(PO<sub>3</sub>)<sub>3</sub> is at  $2\theta \approx 23^{\circ}$ , which is the same position as that of the weak peak of the pseudocubic  $3 \times 3 \times 3$  superlattice. For the samples of x = 0.06 and 0.09, the weak peaks at  $2\theta \approx 23^\circ$  have almost same intensities, implying the formation of the pseudo-cubic  $3 \times 3 \times 3$  superlattice. However, for the samples of *x* = 0.12, the weak peak at  $2\theta \approx 23^{\circ}$  has larger intensity. This may be also a proof of the appearance of Sc(PO<sub>3</sub>)<sub>3</sub> impurity in the x = 0.12 sample. From the above results, it may be believed that the solid solution limit in  $Sn_{1-x}Sc_xP_2O_7$  was at least x = 0.09.

#### 3.2. SEM image

Typical SEM images of external surface of x = 0.06 and 0.12 pellet samples heat-treated at 873 K for 2 h are displayed in Fig. 2. As seen from Fig. 2, the primary particles have sphere-like shapes, and the particle size distributions are about 200–600 nm in diameter. This is similar to the observation for Al<sup>3+</sup>-doped SnP<sub>2</sub>O<sub>7</sub> prepared through



**Fig. 3.** Temperature dependence of electrical conductivity of  $Sn_{1-x}Sc_xP_2O_7$  (x = 0.03, 0.06, 0.09, 0.12) in wet H<sub>2</sub> atmosphere at 323–523 K.

similar heat treatments [13]. It may be seen also from Fig. 2 that some pores occur in pellet samples, which may result in decrease of the densities of the pellet samples and physical leakage of gas through the samples [9]. The relative densities of the pellet samples prepared in this study were estimated to be 70–80%, which are much lower than traditional ceramics, e.g. YSZ, doped CeO<sub>2</sub>, etc. In order to enhance electrical conduction and application performances, it is necessary to increase relative density of Sc-doped SnP<sub>2</sub>O<sub>7</sub> through adopting more effective preparation methods.

#### 3.3. Effect of $Sc^{3+}$ doping level on the conductivities

Fig. 3 shows Arrhenius plots of the total conductivities of the samples in wet H<sub>2</sub> atmosphere from 323 K to 523 K. It is clear that the conductivities increased in the order:  $\sigma$  (x=0.12)< $\sigma$  $(x=0.03) < \sigma$   $(x=0.09) < \sigma$  (x=0.06). The influence of the doping level of Sc<sup>3+</sup> at Sn<sup>4+</sup> sites on the conductivities may be attributed to the effective concentration of oxygen vacancy and impurity phase of  $Sc(PO_3)_3$ . The higher  $Sc^{3+}$  doping level resulted in the higher oxygen vacancy concentration, nevertheless, the concentrations of point defect pairs,  $Sc'_{Sn}V_0^{\bullet\bullet}$ ,  $Sc'_{Sn}V_0^{\bullet\bullet}Sc'_{Sn}$ , and  $Sc'_{Sn}OH_0^{\bullet}$ , also increase at the same time. Considering the opposite two factors above, the effective concentrations of oxygen vacancy may reach its maximum value at x = 0.06 [22]. As a result, the highest conductivity  $2.76 \times 10^{-2} \,\mathrm{S \, cm^{-1}}$  was observed for the sample of x = 0.06under a wet H<sub>2</sub> atmosphere at 473 K. The conductivities of samples of x > 0.06 decreased with the increasing doping levels due to the formation of the secondary  $Sc(PO_3)_3$  phase and lower effective oxygen vacancy concentrations. The effect of the formation of the pseudo-cubic  $3 \times 3 \times 3$  superlattice on the conductivities is still unclear in present stage.

#### 3.4. Concentration cell measurements

For a concentration cell containing hydrogen and water vapor, the observed values (EMF<sub>obs</sub>) of electromotive forces are given by following Nernst equation [19,24,25].

$$\mathsf{EMF}_{\mathsf{obs}} = \frac{RT}{2F} \left\{ -t_{\mathsf{ion}} \ln \left[ \frac{p_{\mathsf{H}_2(\mathsf{II})}}{p_{\mathsf{H}_2(\mathsf{I})}} \right] + t_0 \ln \left[ \frac{p_{\mathsf{H}_2(\mathsf{O}(\mathsf{II})}}{p_{\mathsf{H}_2(\mathsf{O}(\mathsf{I}))}} \right] \right\}$$



**Fig. 4.** EMFs of the hydrogen concentration cell: H<sub>2</sub>, Pt |Sn<sub>1-x</sub>Sc<sub>x</sub>P<sub>2</sub>O<sub>7</sub> (*x*=0.06)| Pt, H<sub>2</sub>-Ar (*p*<sub>H<sub>2</sub></sub> = 1.01 × 10<sup>4</sup> Pa) and EMFs of the water vapor concentration cell: H<sub>2</sub> (*p*<sub>H<sub>2</sub>O</sub> = 2.34 × 10<sup>3</sup> Pa), Pt| Sn<sub>1-x</sub>Sc<sub>x</sub>P<sub>2</sub>O<sub>7</sub> (*x*=0.06)| Pt, H<sub>2</sub> (*p*<sub>H<sub>2</sub>O</sub> = 1.23 × 10<sup>4</sup> Pa). The operating temperatures are from 373 K to 523 K. Dotted line indicates the theoretical value and symbol ● represents the observed value of EMF of the hydrogen concentration cell at each temperature. Solid line indicates the theoretical value and symbol ▼ represents the observed value of EMF of the hydrogen concentration cell at each temperature.

If the  $p_{H_2}$  in both cell compartments is the same, i.e.  $p_{H_2(II)} = p_{H_2(I)}$ , the EMF of water vapor concentration cell gives the transference numbers of oxide-ions ( $t_0$ ). Similarly, the EMF of hydrogen concentration cell with  $p_{H_2O(II)} = p_{H_2O(I)}$  gives the ionic contribution. The ionic transport number  $t_{ion}$  ( $=t_H + t_0$ ) can be obtained by EMF<sub>obs</sub>/EMF<sub>cal</sub>, where EMF<sub>cal</sub> =  $RT/2F \ln[p_{H_2(II)}/p_{H_2(I)}]$ . Electronic transport number  $t_e$  can be given to  $t_e = 1 - (t_H + t_0)$ .

Typical results using sample of x = 0.06 as an electrolyte of gas concentration cells are shown in Fig. 4. In Fig. 4, pure hydrogen and a mixture of H<sub>2</sub>-Ar ( $p_{H_2} = 1.01 \times 10^4$  Pa) are applied as the electrode gases. The electrode with higher pressure of hydrogen gas is negative [27]. Dotted line indicates the theoretical value and symbol  $\bullet$  represents the observed EMF value of the hydrogen concentration cell at each temperature. Obviously, each observed EMF value of hydrogen concentration cell is close to the corresponding theoretical one. The ionic transport number  $t_{ion}$  is almost unity with value of 0.95–0.99, the rest is electronic transport number ( $t_e = 0.01-0.05$ ) in wet hydrogen atmosphere at 373–523 K. Therefore, the total conductivities of the samples in wet hydrogen atmosphere shown in Fig. 3 is almost pure ionic.

The oxide ionic transport number  $(t_0)$  in wet hydrogen atmosphere may also be evaluated from water vapor concentration cell shown in Fig. 4. Solid line indicates the theoretical value and symbol **v**represents the observed EMF value of the water vapor concentration cell at each temperature. The  $t_0$  is calculated to be 0.13–0.20 at 373–523 K, indicating the existence of oxide ionic conduction under low oxygen partial pressure condition. Thus, the protonic transport number ( $t_H = t_{ion} - t_0$ ) is calculated to be 0.79–0.82. In wet hydrogen atmosphere, protonic conduction appears due to defect reactions (1) and (2) [9], resulting in the prevailing protonic conduction in wet hydrogen atmosphere [22].

$$H_2O(g) + V_0^{\bullet\bullet} + O_0^{\times} \Leftrightarrow 2OH_0^{\bullet}$$
(1)

$$H_2O(g) + 2h^{\bullet} + 2O_0^{\times} \Leftrightarrow 2OH_0^{\bullet} + 1/2O_2$$
(2)



Fig. 5. Partial conductivities of charged species in wet hydrogen atmosphere. The operating temperatures are from 373 K to 523 K.

#### 3.5. Partial conductivities of charged species

In wet hydrogen atmosphere, the total conductivity  $\sigma_t = \sigma_H + \sigma_o + \sigma_e$ . The protonic, oxide ionic and electronic conductivities are calculated from  $\sigma_H = \sigma_t \times t_H$ ,  $\sigma_0 = \sigma_t \times t_0$  and  $\sigma_e = \sigma_t \times t_e$ , respectively. The results are shown in Fig. 5. It is clear from Fig. 5 that the total conductivity ( $\sigma_t = 1.96 \times 10^{-2} - 2.76 \times 10^{-2} \, \mathrm{S\,cm^{-1}}$ ) is one to two orders of magnitude higher than the electronic conductivity ( $\sigma_e = 2.74 \times 10^{-4} - 1.21 \times 10^{-3} \, \mathrm{S\,cm^{-1}}$ ). The maximum protonic conductivity is observed to be  $2.24 \times 10^{-2} \, \mathrm{S\,cm^{-1}}$  under wet hydrogen atmosphere at  $873 \, \mathrm{K}$  [28]. The relatively low electronic conductivity in wet hydrogen atmosphere confirmed that this material is a good ionic conductor [9,22].

#### 3.6. Performance of $H_2$ /air fuel cell

Fig. 6 shows the current–voltage–power (*I–V–P*) curves of the H<sub>2</sub>/air fuel cells: wet H<sub>2</sub> ( $p_{H_2O} = 3.2 \times 10^3$  Pa), Pt|Sn<sub>1–x</sub>Sc<sub>x</sub>P<sub>2</sub>O<sub>7</sub> (*x*=0.03, 0.06, 0.09)|Pt, wet Air ( $p_{H_2O} = 3.2 \times 10^3$  Pa). The operating temperature is 423 K, the thickness of electrolytes is 1.7 mm and the flow rates of H<sub>2</sub> and air gases are all 30 ml min<sup>-1</sup> during



**Fig. 6.** *I*–*V*–*P* curves for a hydrogen/air fuel cell using  $Sn_{1-x}Sc_xP_2O_7$  (*x* = 0.03, 0.06, 0.09) as electrolytes at 423 K. Electrolyte thickness: 1.7 mm.

the experiment. The theoretical electromotive force value (EMF) may be calculated according to literature method [22]. As shown in Fig. 6, the open circuit voltages are around 0.97 V, which are lower than the theoretical one (1.216 V at 423 K). The difference may be relevant to the electronic conduction in the samples as shown in Figs. 4 and 5, and relevant to some physical leakage of gas through the electrolyte due to lower relative density of the electrolyte. It can be seen from Fig. 6 that the peak power densities of the cells reached 11.16 mW cm<sup>-2</sup> for x = 0.03, 25.02 mW cm<sup>-2</sup> for x = 0.06 and 14.34 mW cm<sup>-2</sup> for x = 0.09 sample, correspondingly, which is similar to the data reported by Chen et al. for Sn<sub>0.9</sub>In<sub>0.1</sub>P<sub>2</sub>O<sub>7</sub> [19]. The cell performance may be enhanced by doping suitable ions and levels. The results indicated that Sn<sub>1-x</sub>Sc<sub>x</sub>P<sub>2</sub>O<sub>7</sub> is a promising solid electrolyte system for intermediate temperature fuel cells.

#### 4. Conclusions

In this study, a novel series of samples  $Sn_{1-x}Sc_xP_2O_7$  (x = 0.03, 0.06, 0.09, 0.12, 0.15) were prepared. The doping limit of  $Sc^{3+}$ in  $SnP_2O_7$  was at least x = 0.09. It was found that the conductivities increased with various  $\mathrm{Sc}^{3+}$  doping levels in the order:  $\sigma$  $(x=0.12) < \sigma (x=0.03) < \sigma (x=0.09) < \sigma (x=0.06)$ . The highest conductivity was observed to be  $2.76 \times 10^{-2} \,\text{S}\,\text{cm}^{-1}$  for the sample of x = 0.06 under wet H<sub>2</sub> atmosphere at 473 K. The ionic transport numbers ( $t_{ion} = 0.95 - 0.99$ ) are close to unity, and the relatively low electronic conductivity in wet hydrogen atmosphere indicated that this material is a good ionic conductor. The maximum proton conductivity  $(2.24 \times 10^{-2} \,\text{S}\,\text{cm}^{-1})$  at 473 K is higher than that in BaCe<sub>0.85</sub>Y<sub>0.15</sub>O<sub>3- $\alpha$ </sub> (1.04  $\times$  10<sup>-2</sup> S cm<sup>-1</sup>) under wet hydrogen atmosphere at 873 K. The H<sub>2</sub>/air fuel cells using  $Sn_{1-x}Sc_xP_2O_7$ (x=0.03, 0.06, 0.09) as electrolytes (thickness: 1.7 mm) generated the maximum power densities of 11.16 mW cm<sup>-2</sup> for x = 0.03,  $25.02 \text{ mW cm}^{-2}$  for x=0.06 and  $14.34 \text{ mW cm}^{-2}$  for x=0.09 at 423 K, respectively.  $Sn_{1-x}Sc_xP_2O_7$  may be a promising solid electrolyte system for intermediate temperature fuel cells.

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